

TOPOLOGICAL RECURSION WITH HARD EDGES

LEONID CHEKHOV AND PAUL NORBURY

ABSTRACT. We prove a Givental type decomposition for partition functions that arise out of topological recursion applied to spectral curves. Copies of the Kontsevich-Witten KdV tau function arise out of regular spectral curves and copies of the Brezin-Gross-Witten KdV tau function arise out of irregular spectral curves. We present the example of this decomposition for the matrix model with two hard edges and spectral curve $(x^2 - 4)y^2 = 1$.

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1. INTRODUCTION

The partition functions for Gromov-Witten invariants of \mathbb{P}^1 , the Gaussian Hermitian matrix model, the Legendre ensemble, and enumeration of dessins d'enfant, which are formal series in an infinite sequence of variables $\{\hbar, v^{d,\alpha} \mid d \in \mathbb{N}, \alpha \in \{1, 2\}\}$, have in common a decomposition given by a differential operator \hat{R} acting on the product of two species of the Kontsevich-Witten KdV tau function Z^{KW} or of the Brezin-Gross-Witten KdV tau function Z^{BGW} —defined in Section 2.

$$\begin{aligned}
 (1.1) \quad Z^{\text{GW}}(\hbar, \{v^{d,\alpha}\}) &= \hat{R} \cdot \hat{T}_1 \cdot Z^{\text{KW}}(2\hbar, \{\sqrt{2}v^{d,1}\}) Z^{\text{KW}}(-2\hbar, \{i\sqrt{2}v^{d,2}\}) \\
 Z^{\text{GUE}}(\hbar, \{v^{d,\alpha}\}) &= \hat{R} \cdot \hat{T}_2 \cdot Z^{\text{KW}}(2\hbar, \{\sqrt{2}v^{d,1}\}) Z^{\text{KW}}(-2\hbar, \{i\sqrt{2}v^{d,2}\}) \\
 Z^{\text{Leg}}(\hbar, \{v^{d,\alpha}\}) &= \hat{R} \cdot \hat{T}_3 \cdot Z^{\text{BGW}}(2\hbar, \{\sqrt{2}v^{d,1}\}) Z^{\text{BGW}}(-2\hbar, \{i\sqrt{2}v^{d,2}\}) \\
 Z^{\text{Des}}(\hbar, \{v^{d,\alpha}\}) &= \hat{R} \cdot \hat{T}_4 \cdot Z^{\text{BGW}}\left(-\frac{1}{2}\hbar, \left\{\frac{i}{\sqrt{2}}v^{d,1}\right\}\right) Z^{\text{KW}}(32\hbar, \{4\sqrt{2}v^{d,2}\}).
 \end{aligned}$$

The operator \hat{R} , which is the exponential of a quadratic differential operator, is common to all four models, whereas \hat{T}_i are operators of constant translations $v^{d,\alpha} \mapsto v^{d,\alpha} + c_i^{d,\alpha}$ described in Section 4.1. The partition function Z^{GW} stores ancestor Gromov-Witten invariants of \mathbb{P}^1 . Its decomposition in (1.1) is a particular

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case of Givental's decomposition of partition functions of Gromov-Witten invariants of targets X with semi-simple quantum cohomology [20] which applies more generally to partition functions of cohomological field theories. It is usually expressed as a function of variables corresponding to cohomology classes in $H^*(\mathbb{P}^1)$ denoted $\{t^{d,\beta}\}$ which are related to the variables in the decomposition by $v^{d,1} = \frac{1}{\sqrt{2}}(t^{d,1} + t^{d,2})$, $v^{d,2} = \frac{i}{\sqrt{2}}(-t^{d,1} + t^{d,2})$. The partition functions Z^{GUE} and Z^{Leg} store moments of the probability measure $\int_{H_N} \exp V(M) DM$ as asymptotic expansions in $\hbar = 1/N$ for $N \rightarrow \infty$ where H_N consists of $N \times N$ Hermitian matrices. For Z^{GUE} , we use $V(M) = -\text{tr}(M^2)$, for Z^{Des} we use $V(M) = -\text{tr} M$ and we have an infinite wall at zero, and for Z^{Leg} we use $V(M)$ given by an infinite well potential, meaning that one restricts eigenvalues to lie in the interval $[-2, 2]$ and sets $V(M) = 0$ —for the integration over a compact domain we do not need a Gaussian term for convergence. The decomposition of Z^{GUE} is a decomposition of the partition function for a Hermitian matrix model with Gaussian potential proven by the first author in [5] and in fact gives an example of Givental's decomposition via an associated cohomological field theory [3]. The decomposition of Z^{Leg} into two copies of Z^{BGW} is described in detail in this paper as a particular example of the more general result involving copies of both Z^{KW} and Z^{BGW} . An example of mixed Z^{BGW} and Z^{KW} factors is given by Z^{Des} the partition function for enumeration of dessins d'enfant [9]. The conclusion is that the partition functions Z^{KW} and Z^{BGW} are fundamental to a large class of partition functions arising from many areas.

For the two choices of $V(M)$ above, the limit

$$y(x) = \lim_{N \rightarrow \infty} N^2 \int_{H_N} \left\langle \text{tr} \frac{1}{x - M} \right\rangle \exp V(M) DM$$

is a holomorphic function which defines a Riemann surface, known as a spectral curve, given as a double cover of the x -plane $x = z + 1/z$ on which $y(x)dx$ extends to a well-defined differential $r(z)dz$ for $r(z)$ a rational function. One can also associate a Riemann surface to Gromov-Witten invariants of \mathbb{P}^1 via an associated Landau-Ginzburg model. Each of the examples in (1.1) can be formulated in terms of a recursive construction of holomorphic differentials defined on the associated Riemann surface known as *topological recursion*. It is with respect to this formulation of partition functions that we prove a rather general decomposition theorem.

Topological recursion developed by Eynard, Orantin and the first author [6, 7, 17] produces invariants $\omega_{g,n}$ for integers $g \geq 0$ and $n \geq 1$, which we will refer to as *correlators*, of a Riemann surface Σ equipped with two meromorphic functions $x, y : \Sigma \rightarrow \mathbb{C}$ and a bidifferential $B(p_1, p_2)$ for $p_1, p_2 \in \Sigma$. The zeros \mathcal{P}_α of dx must be simple and disjoint from the zeros of dy . We refer to the data $S = (\Sigma, B, x, y)$ as a *spectral curve*. We allow Σ to be (a possibly disconnected) open subset of a compact Riemann surface, in which case S is known as a *local spectral curve*. For integers $g \geq 0$ and $n \geq 1$, the correlator $\omega_{g,n}$ is a multidifferential on Σ or, in other words, a tensor product of meromorphic differentials on Σ^n . It is defined recursively via

$$\omega_{0,1}(p) = -y(p)dx(p), \quad \omega_{0,2}(p_1, p_2) = B(p_1, p_2)$$

which are used to define the kernel in a neighbourhood of $p_2 = \mathcal{P}_\alpha$ for $dx(\mathcal{P}_\alpha) = 0$

$$K(p_1, p_2) = \frac{1}{2} \frac{\int_{\sigma_\alpha(p_2)}^{p_2} B(p, p_1)}{(y(p_2) - y(\sigma_\alpha(p_2)))dx(p_2)}.$$

The point $\sigma_\alpha(p) \in \Sigma$ is defined to be the unique point $\sigma_\alpha(p) \neq p$ close to α such that $x(\sigma_\alpha(p)) = x(p)$ which is well-defined since each zero \mathcal{P}_α of dx is assumed to be simple. For $L = \{2, \dots, n\}$ define

$$(1.2) \quad \omega_{g,n}(p_1, \mathbf{p}_L) = \sum_{\alpha=1}^N \text{Res}_{p=\mathcal{P}_\alpha} K(p_1, p) \left[\omega_{g-1, n+1}(p, \sigma_\alpha(p), \mathbf{p}_L) + \sum_{\substack{g_1+g_2=g \\ L \sqcup J=L}} \omega_{g_1, |I|+1}(p, \mathbf{p}_I) \omega_{g_2, |J|+1}(\sigma_\alpha(p), \mathbf{p}_J) \right]$$

where the outer summation is over the zeros \mathcal{P}_α of dx and the \circ over the inner summation means that we exclude terms that involve ω_1^0 . The recursive definition of $\omega_{g,n}(p_1, \dots, p_n)$ uses only local information around

zeros of dx so a local spectral curve containing the zeros of dx is sufficient. A zero of dx is *regular* if y is analytic there. A spectral curve is regular if y is analytic at all zeros of dx . In this paper we consider *irregular* spectral curves where y may have a simple pole at any zero of dx . The correlators $\omega_{g,n}$ are polynomial in a basis of differentials $v^{k,m} = V^{k,m}(p_i)$ on Σ depending only x and B —defined in (3.1) in Section 3. Define the topological recursion partition function of the spectral curve $S = (\Sigma, B, x, y)$ by

$$Z^S(\hbar, \{v^{d,i}\}) = \exp\left(\sum_{g,n} \frac{\hbar^{g-1}}{n!} \omega_{g,n}(\{v^{d,i}\})\right).$$

The topological recursion partition function of the curve $x = \frac{1}{2}y^2$ (equipped with the Cauchy kernel $B = dy_1 dy_2 / (y_1 - y_2)^2$) is the Kontsevich-Witten KdV tau function $Z = Z_{\text{KW}}$ and we write $\omega_{g,n}^{\text{Airy}}$ for the correlators of this curve known as the Airy curve due to its relation to the differential equation satisfied by the Airy function. Similarly, the topological recursion partition function of the curve $xy^2 = \frac{1}{2}$ yields the Brezin-Gross-Witten KdV tau function $Z = Z_{\text{BGW}}$, defined in Section 2, and we write $\omega_{g,n}^{\text{Bes}}$ due to a relation of the curve with the Bessel equation [10]. The BGW model was first identified with the KdV τ -function in [27].

For a general spectral curve S , it is straightforward to prove that the asymptotic behaviour, or largest polar part, of $\omega_{g,n}$ near each zero \mathcal{P}_α of dx is given by the correlators for the local model of the curve $x = \frac{1}{2}y^2$ and $xy^2 = \frac{1}{2}$, i.e. $\omega_{g,n} \sim \omega_{g,n}^{\text{Airy}}$ near any regular zero of dx and $\omega_{g,n} \sim \omega_{g,n}^{\text{Bes}}$ near any irregular zero of dx where y has a simple pole. What is much deeper is that $\omega_{g,n}$ can be constructed completely from copies of $\omega_{g,n}^{\text{Airy}}$ and $\omega_{g,n}^{\text{Bes}}$. This is described in terms of the partition functions in Theorem 1 below.

In [12] Dunin-Barkowski, Orantin, Shadrin and Spitz proved that the partition function Z^S of a regular spectral curve satisfying a finiteness assumption possesses a decomposition in terms of products of Z^{KW} acted on by differential operators built out of spectral curve data. Furthermore, they showed that this decomposition coincides with a decomposition of Givental [20] for partition functions Z arising out of cohomological field theories. An immediate consequence is that, under some assumptions on the spectral curve, topological recursion produces partition functions Z for cohomological field theories.

The results of [12] require the spectral curve to have regular singularities, i.e. dy must be analytic at the zeros of dx . The main result of this paper is a generalisation of the decomposition theorem to allow irregular singularities.

Theorem 1. *Consider a spectral curve $S = (\Sigma, B, x, y)$ with k irregular zeros of dx at which y has simple poles, and $N - k$ regular zeros. There exist operators \hat{R} , \hat{T} and $\hat{\Delta}$ defined in (4.9) determined explicitly by (Σ, B, x, y) such that*

$$(1.3) \quad Z^S = \hat{R} \hat{T} \hat{\Delta} Z^{\text{BGW}}(\hbar, \{v^{d,1}\}) \cdots Z^{\text{BGW}}(\hbar, \{v^{d,k}\}) Z^{\text{KW}}(\hbar, \{v^{d,k+1}\}) \cdots Z^{\text{KW}}(\hbar, \{v^{d,N}\}).$$

Moreover, \hat{R} depends only on (Σ, B, x) , \hat{T} is a translation operator, and $\hat{\Delta}$ acts by rescaling \hbar and $v^{d,k}$.

Remark 1.1. Theorem 1 generalises the result of [12] even when applied to a regular spectral curve since it contains a translation term \hat{T} that does not appear in the decomposition theorem of [12]. More precisely, [12] considers a restricted class of spectral curves on which the function y is determined by (Σ, B, x) together with the N numbers $dy(\mathcal{P}_\alpha) \in \mathbb{C}$ and corresponds to cohomological field theories with flat identity. In particular [12] does not apply to Weil-Petersson volumes studied by Mirzakhani [28] which gives a fundamental example of a cohomological field theory without flat identity. This is despite the proof in [15, 19] that Weil-Petersson volumes do indeed arise out of topological recursion applied to a spectral curve. Theorem 1 remedies this situation. The translation term in Theorem 1 (for regular spectral curves) corresponds to a translation of cohomological field theories—see [30].

Remark 1.2. In the examples (1.1) the different operators $\hat{\Delta}$ are visible via the differential rescalings of \hbar . The operator \hat{R} is built out of the data (Σ, B, x) —see Section 4. It is the same for all four examples of (1.1) reflecting the fact that the associated spectral curves differ only in the definition of y , i.e. (Σ, B, x) is the same in these four cases.

Remark 1.3. The partition functions arising from Gromov-Witten invariants, cohomological field theories, matrix models and topological recursion possess two sets of natural coordinates—*flat* coordinates and *canonical* coordinates. Theorem 1 is expressed with respect to canonical coordinates. The change of coordinates between canonical and flat is given by an $N \times N$ matrix, i.e. it is linear and independent of the first parameter d . This change of coordinates appears in Givental’s decomposition as a further operator acting on the left of (1.3).

The proof of the formula (1.3) is built up progressively via special cases proven throughout the text. The reader may find some of these simpler versions more digestible. The basic cases of $(\hat{R}, \hat{T}, \hat{\Delta}) = (Id, Id, Id)$ and $N = 1$ appear in (2.2) and (2.3). The case of $(\hat{R}, \hat{T}, \hat{\Delta}) = (Id, Id, \hat{\Delta})$ which gives rise to a topological field theory appears in (3.8). Translations can be best understood via the case $(\hat{R}, \hat{T}, \hat{\Delta}) = (Id, \hat{T}, \hat{\Delta})$ and $N = 1$ given in (4.4).

In Section 2 we recall the definitions of the two KdV tau functions which form the fundamental pieces of the decomposition. In Section 3 we introduce the decomposition without the differential operator \hat{R} via the elementary topological part of the correlators. In Section 4 we prove the decomposition (1.3). We apply the decomposition to the example of the Legendre ensemble in Section 5 and demonstrate an application of the decomposition (1.3) pictorially. In this paper we use the convention $\mathbb{N} = \{0, 1, 2, \dots\}$.

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2. KONTSEVICH-WITTEN AND BREZIN-GROSS-WITTEN TAU FUNCTIONS

The fundamental components $Z^{KW}(\hbar, t_0, t_1, \dots)$ and $Z^{BGW}(\hbar, t_0, t_1, \dots)$ of the factorisation (1.3) are tau functions of the KdV hierarchy. The Kontsevich-Witten tau function Z^{KW} was introduced in [31], and the Brezin-Gross-Witten tau function Z^{BGW} arises out of a unitary matrix mode studied in [4, 22]. A tau function $Z(t_0, t_1, \dots)$ of the KdV hierarchy (equivalently the KP hierarchy in odd times $p_{2k+1} = t_k / (2k + 1)!!$) gives rise to a solution of the KdV hierarchy via $Z = \exp F$, $U = \hbar \frac{\partial^2 F}{\partial t_0^2}$

$$(2.1) \quad U_{t_1} = UU_{t_0} + \frac{\hbar}{12} U_{t_0 t_0 t_0}, \quad U(t_0, 0, 0, \dots) = f(t_0).$$

The first equation in the hierarchy is the KdV equation (2.1), and later equations $U_{t_k} = P_k(U, U_{t_0}, U_{t_0 t_0}, \dots)$ for $k > 1$ determine U uniquely from $U(t_0, 0, 0, \dots)$.

The Kontsevich-Witten tau function Z^{KW} is defined by $U(t_0, 0, 0, \dots) = t_0$ (and is in fact determined uniquely by (2.1) and the string equation $F_{t_0} = \frac{1}{2} t_0^2 + \sum t_{i+1} F_{t_i}$ i.e. the higher equations giving U_{t_k} for $k > 1$ are automatically satisfied). It is famously a generating function for intersection numbers over the moduli space of stable curves equipped with tautological line bundles $L_i \rightarrow \overline{\mathcal{M}}_{g,n}$, $i = 1, \dots, n$ and $\psi_i = c_1(L_i)$.

Theorem 2 (Witten-Kontsevich 1992 [23, 31]).

$$F^{KW}(\hbar, t_0, t_1, \dots) = \sum_{g,n} \hbar^{g-1} \frac{1}{n!} \sum_{\vec{k} \in \mathbb{N}^n} \int_{\overline{\mathcal{M}}_{g,n}} \prod_{i=1}^n \psi_i^{k_i} t_{k_i}$$

Its first few terms are given by

$$F^{KW}(\hbar, t_0, t_1, \dots) = \hbar^{-1} \left(\frac{t_0^3}{3!} + \frac{t_0^3 t_1}{3!} + \frac{t_0^4 t_2}{4!} + \dots \right) + \frac{t_1}{24} + \dots$$

Its dispersionless limit is nontrivial: $\lim_{\hbar \rightarrow 0} U = \frac{t_0}{1-t_1} + \frac{t_0^2 t_2}{2(1-t_1)^3} + \dots$

It arises via topological recursion applied to the Airy curve [18]. For

$$S_{\text{Airy}} = \left\{ x = \frac{1}{2}z^2, y = z, B = \frac{dzdz'}{(z-z')^2} \right\}$$

and $\omega_{g,n}^{\text{Airy}}$ defined by (1.2), we have

$$\omega_{g,n}^{\text{Airy}} = \sum_{\vec{k} \in \mathbb{Z}_+^n} \int_{\mathcal{M}_{g,n}} \prod_{i=1}^n \psi_i^{k_i} (2k_i + 1)!! \frac{dz_i}{z_i^{2k_i+2}}.$$

Out of the correlators $\omega_{g,n}^{\text{Airy}}$, (3.2) builds a partition function which in this case using $V^k(z) = (2k+1)!! \frac{dz}{z^{2k+2}}$ from Definition 3.1 gives

$$(2.2) \quad Z^{S_{\text{Airy}}} = Z^{\text{KW}}.$$

This is the first case of (1.3) where $N = 1$ and $\hat{R} = I = \hat{T} = \hat{\Delta}$.

The Brezin-Gross-Witten solution of the KdV hierarchy is defined by the initial condition

$$U(t_0, 0, 0, \dots) = \frac{\hbar}{8(1-t_0)^2}.$$

The first few terms of its tau function are given by

$$\log Z^{\text{BGW}} = F^{\text{BGW}}(\hbar, t_0, t_1, \dots) = \hbar \left(\frac{1}{8}t_0 + \frac{1}{16}t_0^2 + \frac{1}{24}t_0^3 + \dots \right) + \hbar^2 \left(\frac{3}{128}t_1 + \dots \right)$$

and we see that its dispersionless limit is trivial:

$$\lim_{\hbar \rightarrow 0} U = 0.$$

It arises via topological recursion applied to the Bessel curve [10]

$$S_{\text{Bes}} = \left\{ x = \frac{1}{2}z^2, y = \frac{1}{z}, B = \frac{dzdz'}{(z-z')^2} \right\}$$

as follows. Write

$$\omega_{g,n}^{\text{Bes}} = \sum_{\mu \in \mathbb{Z}_+^n} b_{g,n}(\mu_1, \dots, \mu_n) \prod_{i=1}^n \frac{dz_i}{z_i^{\mu_i+1}}$$

for the correlators of topological recursion applied to the curve S_{Bes} . As above $V^k(z) = (2k+1)!! \frac{dz}{z^{2k+2}}$ and it is proven in [10] that

$$F_g^{S_{\text{Bes}}} = F_g^{\text{BGW}} = \sum_n \frac{1}{n!} \sum_{k \in \mathbb{N}^n} \frac{b_{g,n}(2k_1+1, \dots, 2k_n+1)}{\prod_{i=1}^n (2k_i+1)!!} t_{k_1} \dots t_{k_n}$$

hence

$$(2.3) \quad Z^{S_{\text{Bes}}} = Z^{\text{BGW}}$$

which is again a case of (1.3) for $N = 1$ and $\hat{R} = I = \hat{T} = \hat{\Delta}$.

Remark 2.1. Kontsevich and Soibelman [24] have studied topological recursion via an algebraic structure which they call an *Airy structure* referring to the fact that a regular spectral curve locally resembles the Airy curve. The more general setup of irregular spectral curves that can locally resemble the Bessel curve also fits into the picture of Kontsevich and Soibelman using the technique of abstract topological recursion [2].

3. PARTITION FUNCTION FOR TOPOLOGICAL RECURSION

In this section we define the partition function Z^S built out of the correlators $\omega_{g,n}$ of the spectral curve S . It is a generating function for all $\omega_{g,n}$ with the substitution of variables $v^{d,i}$ for differentials on the curve. We then give a leisurely introduction to the formula (1.3) by considering only part of this formula, obtained by arranging $\hat{R} = Id = \hat{T}$.

Any correlator $\omega_{g,n}(p_1, \dots, p_n)$ has the property that its principal part in any p_i at any zero \mathcal{P}_i of dx is skew-invariant under the local involution defined by dx around \mathcal{P}_i . Eynard [14] defined a collection of auxiliary differentials (defined below) on the curve which span all those meromorphic differentials with principal part at any zero \mathcal{P} of dx skew-invariant under the local involution defined by dx around \mathcal{P} . Hence $\omega_{g,n}$ is a polynomial in these auxiliary differentials.

Definition 3.1. For a Riemann surface Σ equipped with a meromorphic function $x : \Sigma \rightarrow \mathbb{C}$ define the auxiliary differentials on Σ as follows:

$$(3.1) \quad V_0^\alpha(p) = B(\mathcal{P}_\alpha, p), \quad V_{k+1}^\alpha(p) = d \left(\frac{V_k^\alpha(p)}{dx(p)} \right), \quad \alpha = 1, \dots, N, \quad k = 0, 1, 2, \dots$$

where $B(\mathcal{P}_\alpha, p)$ is evaluation at \mathcal{P}_α , a zero of dx . Evaluation of any meromorphic differential ω at a simple zero \mathcal{P} of dx is defined by

$$\omega(\mathcal{P}) := \operatorname{Res}_{p=\mathcal{P}} \frac{\omega(p)}{\sqrt{2(x(p) - x(\mathcal{P}))}}$$

where we choose a branch of $\sqrt{x(p) - x(\mathcal{P})}$ once and for all at each \mathcal{P} to remove the ± 1 ambiguity. The thus defined meromorphic differentials constitute local Krichever–Whitham systems [25] of 1-differentials symmetric with respect to local involutions.

Note that in [14] Eynard defines $V_k^\alpha(p)$ (written $d\tilde{\xi}_{\alpha,k}(p)$) using local coordinates. For each locally defined involution σ_α defined in a neighbourhood of a zero \mathcal{P}_α of dx , we have $V_k^\alpha(p) + V_k^\alpha(\sigma_\alpha(p))$ is analytic at \mathcal{P}_α . The $V_k^\alpha(p)$ form a basis for meromorphic differentials which have principal part skew invariant under each involution σ_α and $\omega_{g,n}$ is a polynomial in them:

$$\omega_{g,n}(p_1, \dots, p_n) = \sum_{\vec{\alpha}, \vec{k}} c_{g, \vec{\alpha}, \vec{k}} \prod_{i=1}^n V_{k_i}^{\alpha_i}(p_i).$$

The partition function of a spectral curve $S = (\Sigma, B, x, y)$ is defined by:

$$(3.2) \quad Z^S(\hbar, \{v^{k,\alpha}\}) = \exp \left(\sum_{g,n} \hbar^{2g} \frac{1}{n!} \omega_{g,n}(p_1, \dots, p_n) \Big|_{\{V_{k_i}^{\alpha_i}(p_i) = v^{k_i, \alpha_i}\}} \right).$$

3.1. Topological field theory and asymptotic behaviour of $\omega_{g,n}$. Given a spectral curve $S = (\Sigma, B, x, y)$, around any regular zero \mathcal{P}_i of dx the pair (Σ, x) resembles the Airy curve $x = \frac{1}{2}z^2$. A consequence of this proven in [17] is that near a branch point, the asymptotic behaviour of $\omega_{g,n}(p_1, \dots, p_n)$ is described by $\omega_{g,n}^{\text{Airy}}(z_1, \dots, z_n)$. More precisely, consider a local variable s in a neighbourhood of a_i chosen so that $x = x(a_i) + \frac{1}{2}s^2$. With respect to this local coordinate $y = y(\mathcal{P}_i) + \eta_i^{1/2}s + \dots$ where

$$(3.3) \quad \eta_i = dy(\mathcal{P}_i)^2 = \operatorname{Res}_{z=\mathcal{P}_i} \frac{dydy}{dx}.$$

The dominant asymptotic term as $s \rightarrow 0$ is

$$(3.4) \quad \omega_{g,n} = s^{6-6g-3n} dy(\mathcal{P}_i)^{2-2g-n} \omega_{g,n}^{\text{Airy}} + O(s^{5-6g-3n}).$$

The analogous behaviour near an irregular zero \mathcal{P}_i of dx where the spectral curve resembles the Bessel curve $xy^2 = \frac{1}{2}$ also holds. With respect to the local variable s such that $x = x(\mathcal{P}_i) + \frac{1}{2}s^2$ we have $y = \eta_i^{1/2}s^{-1} + \dots$ for

$$(3.5) \quad \eta_i = (ydx)(\mathcal{P}_i)^2 = \operatorname{Res}_{z=\mathcal{P}_i} y^2 dx.$$

The dominant asymptotic term as $s \rightarrow 0$ is

$$(3.6) \quad \omega_{g,n} = s^{-2g-2n} dy(\mathcal{P}_i)^{-g-n} \omega_{g,n}^{\text{Bessel}} + O(s^{-1-2g-2n}).$$

Collect the top order pole parts of $\omega_{g,n}(p_1, \dots, p_n)$ at each \mathcal{P}_i into a single correlator $\omega_{g,n}^{\text{top}}(p_1, \dots, p_n)$. An alternative definition of $\omega_{g,n}^{\text{top}}(p_1, \dots, p_n)$ uses a local spectral curve S_0 built out of S . Given a spectral curve $S = (\Sigma, B, x, y)$, define the local spectral curve $S_0 = (\Sigma, B_0, x, y)$ with B_0 the trivial Bergmann kernel—it is given in a local coordinate s around any zero \mathcal{P}_i of dx defined by $x = x(\mathcal{P}_i) + \frac{1}{2}s^2$ by $B_0(p, p') = \frac{ds(p)ds(p')}{(s(p)-s(p'))^2}$. The correlators $\omega_{g,n}^{\text{top}}(p_1, \dots, p_n)$ of S_0 consist of the top order pole parts of $\omega_{g,n}^{\text{top}}(p_1, \dots, p_n)$ at each zero of dx .

In preparation for studying the full formula (1.3), we restate (3.4) and (3.6) by (3.8) below in terms of the partition functions Z_{KW} and Z_{BGW} by replacing the \hat{R} operator in (1.3) with the identity operator.

Recall that a two-dimensional topological field theory (2D TFT) is a vector space H and a sequence of symmetric linear maps

$$I_{g,n} : H^{\otimes n} \rightarrow \mathbb{C}$$

for integers $g \geq 0$ and $n > 0$ satisfying the following conditions. The map $I_{0,2}$ is a non-degenerate bilinear form on H i.e. it defines a metric η , with dual bivector $\Delta = \eta^{\alpha\beta} e_\alpha \otimes e_\beta$ (defined with respect to a basis $\{e_\alpha\}$ of H). The map $I_{0,3}$ together with $\eta = I_{0,2}$ defines a product \cdot on H via

$$\eta(v_1 \cdot v_2, v_3) = I_{0,3}(v_1, v_2, v_3)$$

with identity $\mathbb{1}$ given by the dual of $I_{0,1} = \mathbb{1}^* = \eta(\mathbb{1}, \cdot)$. It satisfies $I_{g,n+1}(\mathbb{1} \otimes v_1 \otimes \dots \otimes v_n) = I_{g,n}(v_1 \otimes \dots \otimes v_n)$ and the gluing conditions

$$I_{g,n}(v_1 \otimes \dots \otimes v_n) = I_{g-1,n+2}(v_1 \otimes \dots \otimes v_n \otimes \Delta) = I_{g_1,|I|+1} \otimes I_{g_2,|J|+1} \left(\bigotimes_{i \in I} v_i \otimes \Delta \otimes \bigotimes_{j \in J} v_j \right)$$

for $g = g_1 + g_2$ and $I \sqcup J = \{1, \dots, n\}$.

Via the natural isomorphism $H^0(\overline{\mathcal{M}}_{g,n}) \cong \mathbb{C}$ we consider $I_{g,n} : H^{\otimes n} \rightarrow H^0(\overline{\mathcal{M}}_{g,n})$ and integrate these classes next to Chern classes of the tautological line bundles \mathcal{L}_i on the moduli space of stable curves $\overline{\mathcal{M}}_{g,n}$. With respect to a basis $\{e_{v_i}\}$ of H , we define the partition function:

$$(3.7) \quad Z(\hbar, \{v^{d,j}\}) = \exp \sum_{g,n} \hbar^{g-1} \frac{1}{n!} \int_{\overline{\mathcal{M}}_{g,n}} I_{g,n}(e_{v_1}, \dots, e_{v_n}) \cdot \prod_{j=1}^n c_1(\mathcal{L}_j)^{d_j} v^{d_j v_j}.$$

The partition function (3.7) for the trivial dimension 1 TFT, where $I_{g,n}(\mathbb{1}^{\otimes n}) = 1$, stores intersection numbers of ψ classes on $\overline{\mathcal{M}}_{g,n}$ and hence $Z(\hbar, \{v^{d,1}\}) = Z^{\text{KW}}(\hbar, \{v^{d,1}\})$, the Kontsevich-Witten partition function, as described in Section 2.

A 2D TFT is known as *semisimple* if its associated Frobenius algebra (H, η, \cdot) is *semisimple*, i.e.

$$H \cong \mathbb{C} \oplus \mathbb{C} \oplus \dots \oplus \mathbb{C}, \quad \langle e_i, e_j \rangle = \delta_{ij} \eta_i, \quad e_i \cdot e_j = \delta_{ij} e_i$$

for some $\eta_i \in \mathbb{C} \setminus \{0\}$, $i = 1, \dots, N$ where $e_i^{(j)} = \delta_{ij}$ is the standard basis. For a semisimple 2D TFT, the $I_{g,n}$ decompose into a sum of 1-dimensional TFTs and hence are extremely simple. A 1-dimensional TFT depends on a single complex number $I_{0,1}(\mathbb{1}) = \eta \in \mathbb{C}$ which determines $I_{g,n}(\mathbb{1}^{\otimes n}) = \eta^{1-g}$. Its partition function (3.7) is simply $Z^{\text{KW}}(\eta^{-1} \hbar, \{u^{d,1}\})$ since the η^{1-g} is naturally absorbed by the \hbar^{g-1} . The coordinate $u^{d,1}$ corresponds to the unit vector $\mathbb{1}$, and we instead use $v^{d,1} = \eta^{1/2} u^{d,1}$ corresponding to an orthonormal

basis, so $Z^{\text{KW}}(\eta^{-1}\hbar, \{u^{d,1}\}) = Z^{\text{KW}}(\eta^{-1}\hbar, \{\eta^{-\frac{1}{2}}v^{d,1}\})$. Hence the partition function of a semisimple 2D TFT, where $I_{g,n}$ vanishes on mixed monomials in the e_i and $I_{g,n}(e_i^{\otimes n}) = \eta_i^{1-g}$, is given by a product

$$Z(\hbar, \{v^{d,j}\}) = Z_{\text{KW}}(\eta_1^{-1}\hbar, \{\eta_1^{-\frac{1}{2}}v^{d,1}\}) \cdots Z_{\text{KW}}(\eta_N^{-1}\hbar, \{\eta_N^{-\frac{1}{2}}v^{d,N}\}).$$

For storing the numbers $I_{g,n}(e_{v_1}, \dots, e_{v_n})$ this may seem more complicated than necessary, however it brings useful insight to the formula (1.3), and is precisely necessary to store the highest order parts $\omega_{g,n}^{\text{top}}(p_1, \dots, p_n)$ of the correlators $\omega_{g,n}$ of a spectral curve $S = (\Sigma, B, x, y)$. A generalisation of this idea couples a TFT to $Z^{\text{BGW}}(\hbar, \{v^{d,1}\})$ via $Z^{\text{BGW}}(\eta^{-1}\hbar, \{\eta^{-\frac{1}{2}}u^{d,1}\})$. Thus

$$(3.8) \quad Z^{S_0}(\hbar, \{v^{m,i}\}) = \hat{\Delta} Z_{\text{BGW}}(\hbar, \{v^{m,1}\}) \cdots Z_{\text{BGW}}(\hbar, \{v^{m,k}\}) Z_{\text{KW}}(\hbar, \{v^{m,k+1}\}) \cdots Z_{\text{KW}}(\hbar, \{v^{m,N}\})$$

where S , hence S_0 , has k irregular zeros of dx at which y has simple poles, and $N - k$ regular zeros, $\hat{\Delta}(\hbar) = \eta_i^{-1}\hbar$ in the i th factor and $\hat{\Delta}(v^{m,k}) = \eta_k^{-\frac{1}{2}}v^{m,k}$.

The expressions (3.4) and (3.6), hence also (3.8), depend only on the local behaviour of x and y in a neighbourhood of \mathcal{P}_i , and are independent of B . The formula (1.3) strengthens this result to show that all lower asymptotic terms of $\omega_{g,n}$, hence $\omega_{g,n}$ itself, can also be obtained from $\omega_{g,n}^{\text{Airy}}$ and $\omega_{g,n}^{\text{Bessel}}$ by also using B .

4. GIVENTAL DECOMPOSITION

4.1. Translations. The translation term \hat{T} in (1.3) translates the arguments $v^{d,\alpha}$ in the tau functions:

$$Z^{\text{KW}}(\hbar, \{v^{d,\alpha}\}) \mapsto Z^{\text{KW}}(\hbar, \{v^{d,\alpha} + c^{d,\alpha}\}), \quad Z^{\text{BGW}}(\hbar, \{v^{d,\beta}\}) \mapsto Z^{\text{BGW}}(\hbar, \{v^{d,\beta} + c^{d,\beta}\}).$$

Translations arise from local expansions of y near each zero \mathcal{P}_α of dx . To study these translations, we restrict to the case where dx has a single zero. Let $x = \frac{1}{2}z^2$, $y = \sum_{k=-1}^{\infty} y_k z^k$. This is a deformation of the Bessel curve, or a deformation of the Airy curve when $y_{-1} = 0$ and $y_1 \neq 0$. General deformations of the Airy curve were studied in [12, 15]. Propositions 4.1 and 4.2 below show that the correlators of deformations of the Airy and Bessel curves depend linearly on the coefficients $a_{g,n}$ and $b_{g,n}$ of the correlators $\omega_{g,n}^{\text{Airy}}$ and $\omega_{g,n}^{\text{Bes}}$ defined by

$$(4.1) \quad \omega_{g,n}^{\text{Airy}} = \sum_{\mu \in \mathbb{Z}_+^n} a_{g,n}(\mu) \prod_{i=1}^n \frac{dz_i}{z_i^{\mu_i+1}}, \quad \omega_{g,n}^{\text{Bes}} = \sum_{\mu \in \mathbb{Z}_+^n} b_{g,n}(\mu) \prod_{i=1}^n \frac{dz_i}{z_i^{\mu_i+1}}$$

and homogeneously in the coefficients y_k of y . We will begin with the statement of the known result for deformations of the Airy curve.

Proposition 4.1 ([12, 14, 15]). For $x = \frac{1}{2}z^2$, $y = \sum_{k=1}^{\infty} y_k z^k$, $B = \frac{dzdz'}{(z-z')^2}$

$$(4.2) \quad \omega_{g,n}(z_1, \dots, z_n) = y_1^{2-2g-n} \sum_{m=0}^{g-1} \sum_{\vec{d}, \vec{\alpha}} \prod_{i=1}^n \frac{dz_i}{z_i^{d_i+1}} \frac{(-1)^m}{m!} \prod_{k=1}^m \frac{y_{\alpha_k-2}}{y_1} \frac{1}{\alpha_k} a_{g,n+m}(\vec{d}, \vec{\alpha})$$

where $\vec{d} = (d_1, \dots, d_n)$, $\vec{\alpha} = (\alpha_1, \dots, \alpha_m)$ with $\alpha_k > 3$, $a_{g,n}$ is defined in (4.1) and the product $\prod_{k=1}^m (\cdot)$ is 1 when $m = 0$.

Note that $a_{g,n+m}(\vec{d}, \vec{\alpha}) = 0$ when $|\vec{d}| + |\vec{\alpha}| > 3g - 3 + n$ so the sum is finite.

The statements of Proposition 4.1 in the literature, such as Lemma 3.5 in [12], use intersection numbers on the moduli space of stable curves related via $a_{g,n}(2m_1 + 1, \dots, 2m_n + 1) = \langle \prod_{k=1}^n (2m_k + 1)!! \tau_{m_k} \rangle_{g,n}$.

Proposition 4.2. For $x = \frac{1}{2}z^2$, $y = \sum_{k=-1}^{\infty} y_k z^k$, $B = \frac{dzdz'}{(z-z')^2}$

$$(4.3) \quad \omega_{g,n}(z_1, \dots, z_n) = y_{-1}^{2-2g-n} \sum_{m=0}^{g-1} \sum_{\vec{d}, \vec{\alpha}} \prod_{i=1}^n \frac{dz_i}{z_i^{d_i+1}} \frac{(-1)^m}{m!} \prod_{k=1}^m \frac{y_{\alpha_k-2}}{y_{-1}} \frac{1}{\alpha_k} b_{g,n+m}(\vec{d}, \vec{\alpha})$$

where $\vec{d} = (d_1, \dots, d_n)$, $\vec{\alpha} = (\alpha_1, \dots, \alpha_m)$ with $\alpha_k > 1$, $b_{g,n}$ is defined in (4.1) and the product $\prod_{k=1}^m (\cdot)$ is 1 when $m = 0$.

Note that $b_{g,n+m}(\vec{d}, \vec{\alpha}) = 0$ when $|\vec{d}| + |\vec{\alpha}| \neq 2g - 2 + n$ so the sum is finite.

Proof. The dependence of the kernel $K(z_1, z)$ on y_{2k-1} in the recursion (1.2) is given by

$$\begin{aligned} K(z_1, z) &= \frac{1}{2} \frac{\int_{-z}^z B(z', z_1)}{(y(z) - y(-z))dx(z)} = \frac{dz_1}{2(z_1^2 - z^2) \sum_{k=0}^{\infty} y_{2k-1} z^{2k-1} dz} \\ &= \frac{1}{y_{-1}} \frac{zdz_1}{2(z_1^2 - z^2) \left(1 + \sum_{k=1}^{\infty} \frac{y_{2k-1}}{y_{-1}} z^{2k}\right) dz} = \frac{1}{y_{-1}} \frac{zdz_1}{2(z_1^2 - z^2) dz} \sum_{m=0}^{\infty} (-1)^m \left(\sum_{k=1}^{\infty} \frac{y_{2k-1}}{y_{-1}} z^{2k} \right)^m. \end{aligned}$$

Thus each summand contributes a term of homogeneous degree m in the y_{2k-1} for $k \geq 1$ to $K(z_1, z)$. Thus we can write $\omega_{g,n}$ as

$$\begin{aligned} \omega_{g,n}(z_1, z_2, \dots, z_n) &= \operatorname{Res}_{z=0} K(z_1, z) \left[\omega_{g-1, n+1}(z, -z, z_2, \dots, z_n) + \sum_{\substack{g_1+g_2=g \\ I \sqcup J = \{2, \dots, n\}}} \omega_{g_1, |I|+1}(z, z_I) \omega_{g_2, |J|+1}(-z, z_J) \right] \\ &= y_{-1}^{2-2g-n} \sum_{m=0}^{g-1} \sum_{\vec{d}, \vec{\alpha}} \prod_{i=1}^n \frac{dz_i}{z_i^{d_i+1}} (-1)^m \prod_{k=1}^m \frac{y_{\alpha_k}}{y_{-1}} C_{g,n+m}(\vec{d}, \vec{\alpha}) \end{aligned}$$

for some constants $C_{g,n+m}(\vec{d}, \vec{\alpha})$ which we will show coincide with coefficients $\frac{1}{m!} b_{g,n+m}(\vec{d}, \vec{\alpha} + 2)$ of $\omega_{g,n}^{\text{Bes}}$ (for $\vec{\alpha} + 2 = (\alpha_1 + 2, \dots, \alpha_k + 2)$). The dependence on only odd α_k is realised by the same property of $b_{g,n+m}(\vec{d}, \vec{\alpha})$. It is easy to prove by induction that $\omega_{g,n}$ is a polynomial in the y_k for odd positive k .

A variational formula for a family of spectral curves depending on a parameter was proven in [8, 17]. When the variation of ydx is given by integration of $B(p_1, p_2)$ around a generalised cycle

$$\frac{\partial}{\partial y_k} (-ydx) = -z^{k+1} dz = -\frac{1}{k+2} d(z^{k+2}) = -\operatorname{Res}_{z_0=z} \frac{1}{k+2} z_0^{k+2} B(z_0, z) = \operatorname{Res}_{z_0=\infty} \frac{1}{k+2} z_0^{k+2} B(z_0, z)$$

determines the variation of $\omega_{g,n}$ to be integration of $\omega_{g,n+1}$ around the same generalised cycle

$$\frac{\partial}{\partial y_k} \omega_{g,n}(z_1, \dots, z_n) = \operatorname{Res}_{z_0=\infty} \frac{1}{k+2} z_0^{k+2} \omega_{g,n+1}(z_0, z_1, \dots, z_n).$$

Write $[y_i y_j \dots y_k] f$ for the coefficient of $y_i y_j \dots y_k$ in a polynomial f of the variables y_i .

$$\begin{aligned} \left[\prod_{k=1}^m y_{\alpha_k} \right] \omega_{g,n} &= \frac{1}{m!} \frac{\partial^{|\alpha|}}{\partial y_{\alpha_1} \dots \partial y_{\alpha_m}} \omega_{g,n} \Big|_{y_i=0} = \frac{1}{m!} \operatorname{Res}_{z_{n+1}=\infty} \dots \operatorname{Res}_{z_{n+m}=\infty} \prod_{k=1}^m \frac{1}{\alpha_k + 2} z_{n+k}^{\alpha_k+2} \omega_{g,n+m}(z_1, \dots, z_{n+m}) \Big|_{y_i=0} \\ &= \frac{1}{m!} \operatorname{Res}_{z_{n+1}=\infty} \dots \operatorname{Res}_{z_{n+m}=\infty} \prod_{k=1}^m \frac{1}{\alpha_k + 2} z_{n+k}^{\alpha_k+2} \omega_{g,n+m}^{\text{Bes}}(z_1, \dots, z_{n+m}) \\ &= \frac{(-1)^m}{m!} \prod_{k=1}^m \frac{1}{\alpha_k + 2} b_{g,n}(d_1, \dots, d_n, \alpha_1 + 2, \dots, \alpha_m + 2) \end{aligned}$$

since $\omega_{g,n}^{\text{Bes}} = \sum_{\mu \in \mathbb{Z}_+^n} b_{g,n}(\mu) \prod_{i=1}^n \frac{dz_i}{z_i^{\mu_i+1}}$. By shifting each α_k to $\alpha_k - 2$ we get (4.3) as required. \square

It may be helpful to read (4.3) via small genus examples:

$$\begin{aligned} \omega_{1,n} &= y_{-1}^{-n} \prod_{i=1}^n \frac{dz_i}{z_i^2} \cdot \frac{1}{8} (n-1)! = y_{-1}^{-n} \prod_{i=1}^n \frac{dz_i}{z_i^2} \cdot b_{1,n}(1, 1, \dots, 1) \\ \omega_{2,n} &= y_{-1}^{-n-2} \prod_{i=1}^n \frac{dz_i}{z_i^2} \left(\sum_{j=1}^n \frac{1}{z_j^2} \frac{9}{256} (n+1)! - \frac{y_1}{y_{-1}} \frac{3}{256} (n+2)! \right) \\ &= y_{-1}^{-n-2} \prod_{i=1}^n \frac{dz_i}{z_i^2} \left(\sum_{j=1}^n \frac{1}{z_j^2} b_{2,n}(3, 1, \dots, 1) - \frac{y_1}{y_{-1}} \frac{1}{3} b_{2,n+1}(1, \dots, 1, 3) \right) \\ \omega_{3,n} &= y_{-1}^{-n-4} \prod_{i=1}^n \frac{dz_i}{z_i^2} \left(\sum_{j=1}^n \frac{1}{z_j^4} \frac{75}{8192} (n+3)! + \sum_{i,j} \frac{1}{z_i^2 z_j^2} \frac{189}{20480} (n+3)! \right. \\ &\quad \left. - \left[\sum_j \frac{1}{z_j^2} \frac{y_1}{y_{-1}} \frac{63}{20480} (n+4)! + \frac{y_3}{y_{-1}} \frac{15}{8192} (n+4)! \right] + \frac{y_1^2}{y_{-1}^2} \frac{21}{40960} (n+5)! \right) \\ &= y_{-1}^{-n-4} \prod_{i=1}^n \frac{dz_i}{z_i^2} \left(\sum_{j=1}^n \frac{1}{z_j^4} b_{3,n}(5, 1, \dots, 1) + \sum_{i,j} \frac{1}{z_i^2 z_j^2} b_{3,n}(3, 3, 1, \dots, 1) \right. \\ &\quad \left. + (-1) \left[\sum_j \frac{1}{z_j^2} \frac{y_1}{y_{-1}} \frac{1}{3} b_{3,n+1}(3, 1, \dots, 1, 3) + \frac{y_3}{y_{-1}} \frac{1}{5} b_{3,n+1}(1, 1, \dots, 1, 5) \right] + \frac{(-1)^2}{2!} \frac{y_1^2}{y_{-1}^2} \frac{1}{32} b_{3,n+2}(1, \dots, 1, 3, 3) \right) \end{aligned}$$

Note that we place the arguments of the symmetric function $b_{g,n}$ at each end to emphasise their origin.

The following corollary is a special case of Theorem 1 and in fact is used in the proof of Theorem 1.

Corollary 4.3. *The partition function of the spectral curve*

$$S = (\Sigma, B, x, y) = (\mathbb{P}^1, \frac{dzdz'}{(z-z')^2}, \frac{1}{2}z^2, \sum_{k=-1}^{\infty} y_k z^k)$$

is obtained via translation of the appropriate tau function:

$$Z^S = \begin{cases} Z^{\text{BGW}}(y_{-1}^{-2} \hbar, \{y_{-1}^{-1} v^{d,1} + d!! \frac{y_d}{y_{-1}}\}), & y_{-1} \neq 0 \\ Z^{\text{KW}}(y_1^{-2} \hbar, \{y_1^{-1} v^{d,1} + d!! \frac{y_d}{y_1}\}), & y_{-1} = 0, y_1 \neq 0. \end{cases}$$

or equivalently

$$(4.4) \quad Z^S = \hat{T} \hat{\Delta} Z^{\text{BGW}}(\hbar, \{v^{d,1}\}) \quad \text{or} \quad \hat{T} \hat{\Delta} Z^{\text{KW}}(\hbar, \{v^{d,1}\})$$

for

$$\hat{T} = \exp \left(\sum_{d \text{ odd}} d!! \frac{y_d}{y_{\min}} \partial_{v^{d,1}} \right), \quad \hat{\Delta}(\hbar, \{v^{d,1}\}) = (y_{\min}^{-2} \hbar, \{y_{\min}^{-1} v^{d,1}\}) \quad y_{\min} = y_{-1} \text{ or } y_1.$$

Proof. Propositions 4.1 and 4.3 can be expressed graphically by adding extra leaves, so-called *dilaton leaves*, to each vertex. The weight of a dilaton leaf with label k is $(2k-1)!! \frac{\hbar_{2k-1}}{\hbar_1}$ at a regular vertex and $(2k-1)!! \frac{\hbar_{2k-1}}{\hbar_{-1}}$ at an irregular vertex. But this also the graphical realisation of translation given by the exponential of a constant vector field. \square

Remark 4.4. The regular case of Corollary 4.3 is a consequence of work of Manin and Zograf [26] and Eynard [15] as follows. Consider

$$F^{\text{KW}}(\hbar, t_0, t_1, \dots) = \sum_{g,n} \hbar^{g-1} \frac{1}{n!} \sum_{\vec{k} \in \mathbb{N}^n} \int_{\mathcal{M}_{g,n}} \prod_{i=1}^n \psi_i^{k_i} t_{k_i}$$

and define a generating function for higher Weil-Petersson volumes

$$F^\kappa(\hbar, \vec{t}, \vec{s}) = \sum_{g,n} \hbar^{g-1} \frac{1}{n!} \sum_{\vec{k} \in \mathbb{N}^n} \int_{\mathcal{M}_{g,n}} \prod_{i=1}^n \psi_i^{k_i} t_{k_i} \prod_{j=1}^{\infty} \kappa_j^{m_j} \frac{s_j^{m_j}}{m_j!}.$$

Manin and Zograf [26] proved that F^κ is a translation of F^{KW}

$$F^\kappa(\hbar, \vec{t}, \vec{s}) = F^{\text{KW}}(\hbar, t_0, t_1, t_2 + p_1(\vec{s}), \dots, t_k + p_{k-1}(\vec{s}), \dots)$$

where the p_j are the Schur polynomials defined by

$$1 - \exp\left(-\sum_{i=1}^{\infty} s_i z^i\right) = \sum_{j=1}^{\infty} p_j(s_1, \dots, s_j) z^j.$$

Eynard [15] proved the same relation between higher Weil-Petersson volumes and topological recursion.

Associate to $y(z) = \sum_{k=1}^{\infty} y_k z^k$ (its Laplace transform)

$$\mathcal{L}(f)(z) = \frac{1}{y_1} \sum_{k=1}^{\infty} (2k+1)!! y_{2k+1} z^k = \sum_{j=1}^{\infty} p_j(s_1, \dots, s_j) z^j.$$

Then for $S = (\frac{1}{2}z^2, y(z), \frac{dzdz'}{(z-z')^2})$ he proves $F^S = F^\kappa(\hbar, \vec{t}, \vec{s})$. In other words,

$$F^S = F^{\text{KW}}(\hbar, t_0, t_1, t_2 + 3!! \frac{y_3}{y_1}, \dots, t_{k+1} + (2k+1)!! \frac{y_{2k+1}}{y_1}, \dots)$$

which is the regular case of Corollary 4.3.

4.2. Graphical expansion. In this section we generalise the weighted graphical expansion of the correlators $\omega_{g,n}$ for regular spectral curves proven in [12, 14] to allow for irregular spectral curves.

Given a set $\{1, \dots, N\}$ which will correspond to the zeros of dx on a spectral curve consider the following set of decorated graphs.

Definition 4.5. For a graph γ denote by

$$V(\gamma), \quad E(\gamma), \quad H(\gamma), \quad L(\gamma) = L^*(\gamma) \sqcup L^\bullet(\gamma)$$

its set of vertices, edges, half-edges and leaves. The disjoint splitting of $L(\gamma)$ into ordinary leaves, L^* , and dilaton leaves, L^\bullet , is part of the structure on γ . The set of half-edges consists of leaves and oriented edges so there is an injective map $L(\gamma) \rightarrow H(\gamma)$ and a multiply-defined map $E(\gamma) \rightarrow H(\gamma)$ denoted by $E(\gamma) \ni e \mapsto \{e^+, e^-\} \subset H(\gamma)$. The map sending a half-edge to its vertex is given by $v : H(\gamma) \rightarrow V(\gamma)$. Decorate γ by functions:

$$\begin{aligned} g &: V(\gamma) \rightarrow \mathbb{N} \\ \alpha &: V(\gamma) \rightarrow \{1, \dots, N\} \\ p &: L^*(\gamma) \xrightarrow{\cong} \{p_1, p_2, \dots, p_n\} \subset \Sigma \\ k &: H(\gamma) \rightarrow \mathbb{N} \end{aligned}$$

such that $k|_{L^\bullet(\gamma)} > 1$ and $n = |L^*(\gamma)|$. We write $g_v = g(v)$, $\alpha_v = \alpha(v)$, $\alpha_\ell = \alpha(v(\ell))$, $p_\ell = p(\ell)$, $k_\ell = k(\ell)$. The genus of γ is $g(\gamma) = b_1(\gamma) + \sum_{v \in V(\gamma)} g(v)$ and γ is stable if any vertex labeled by $g = 0$ is of valency ≥ 3 . We

write n_v for the valency of the vertex v . Define $\Gamma_{g,n}$ to be the set of all stable connected genus g decorated graphs with n ordinary leaves.

We now express the correlators of a spectral curve as a sum over decorated graphs, with vertices weighted by coefficients of $\omega_{g,n}^{\text{Airy}}$ or $\omega_{g,n}^{\text{Bes}}$, edges weighted by coefficients of local expansions of B , ordinary leaves weighted by differentials determined by (Σ, B, x) , and dilaton leaves weighted by coefficients of local expansions of y . We follow the exposition in [12] for the regular case and generalise it to the irregular case.

A spectral curve $S = (\Sigma, B, x, y)$ defines a disjoint splitting $V(\gamma) = V^{\text{reg}}(\gamma) \sqcup V^{\text{irreg}}(\gamma)$ of $V(\gamma)$ into regular and irregular vertices. Label the zeros of dx by $\mathcal{P}_1, \dots, \mathcal{P}_N$ and define a vertex v of γ to be irregular if y has a pole at $\mathcal{P}_\alpha(v)$. The spectral curve S also defines weights on any decorated graph $\gamma \in \Gamma_{g,n}$ which will be used to produce correlators $\omega_{g,n}$ of S as weighted sums over all decorations on graphs of type (g, n) . The weights are defined as follows.

Definition 4.6. *Vertex weights.*

$$W(v) = \begin{cases} y_{1,\alpha_v}^{2-2g_v-n_v} a_{g_v,n_v}(\{k_h \mid h \in H(v)\}), & v \in V^{\text{reg}} \\ y_{-1,\alpha_v}^{2-2g_v-n_v} b_{g_v,n_v}(\{k_h \mid h \in H(v)\}), & v \in V^{\text{irreg}} \end{cases}$$

where $y_{1,\alpha_v} = dy(\mathcal{P}_{\alpha_v})$, $y_{-1,\alpha_v} = (ydx)(\mathcal{P}_{\alpha_v})$ and $a_{g,n}(k_1, \dots, k_n)$ and $b_{g,n}(k_1, \dots, k_n)$ —defined in (4.1)—are symmetric functions of k_i so it makes sense to take in a set of cardinality n .

Edge weights.

$$W(e) = (2k_{e^+} - 1)!!(2k_{e^-} - 1)!! B_{2k_{e^+}, 2k_{e^-}}^{\alpha_{e^+}, \alpha_{e^-}}$$

where with respect to the local coordinates z, z' defined by $x = \frac{1}{2}z^2 + x(\mathcal{P}_\alpha)$ and $x = \frac{1}{2}z'^2 + x(\mathcal{P}_{\alpha'})$

$$B^{\alpha, \alpha'}(z, z') = \delta_{\alpha, \alpha'} \frac{dzdz'}{(z - z')^2} + \sum B_{m, m'}^{\alpha, \alpha'} z^m z'^{m'} dzdz'.$$

Here $\{e^+, e^-\}$ are different orientations of the edge e , and by the symmetry $B_{m, m'}^{\alpha, \alpha'} = B_{m', m}^{\alpha', \alpha}$ the weight $B_{2k_{e^+}, 2k_{e^-}}^{\alpha_{e^+}, \alpha_{e^-}}$ depends only on the (unoriented) edge e .

Ordinary leaf weights

$$W(\ell) = V_{k_\ell}^{\alpha_\ell}(p_\ell)$$

The dependence of $\omega_{g,n}$ on $p_i \in \Sigma$ occurs via $V_k^\alpha(p)$ which is defined in (3.1).

Dilaton leaf weights

$$W(\lambda) = \begin{cases} \frac{1}{y_{1,\alpha_\lambda}} (2k_\lambda + 1)!! y_{2k_\lambda+1, \alpha_\lambda}, & v_\lambda \in V^{\text{reg}} \\ \frac{1}{y_{-1,\alpha_\lambda}} (2k_\lambda + 1)!! y_{2k_\lambda+1, \alpha_\lambda}, & v_\lambda \in V^{\text{irreg}} \end{cases}$$

where $y_{2k+1, \alpha}$ are the odd coefficients of the local expansion $y = \sum y_{k, \alpha} z^k$ with respect to the local coordinate z defined by $x = \frac{1}{2}z^2 + x(\mathcal{P}_\alpha)$.

The recursive structure (1.2) can be encoded in graphs [17] with edges decorated by $\omega_{0,1}(p)$, $\omega_{0,2}(p, p')$ and $K(p, p')$ obtained by applying (1.2) repeatedly which corresponds to the construction of a surface of type (g, n) via recursively attaching $2g - 2 + n$ pairs of pants. This enables one to express $\omega_{g,n}$ as a weighted sum over graphs with $2g - 2 + n$ vertices.

Theorem 3. *For a spectral curve (Σ, B, x, y) and $2g - 2 + n > 0$,*

$$\omega_{g,n}(p_1, \dots, p_n) = \frac{1}{n!} \sum_{\gamma \in \Gamma_{g,n}} \prod_{v \in V(\gamma)} W(v) \prod_{e \in E(\gamma)} W(e) \prod_{\ell \in L^*(\gamma)} W(\ell) \prod_{\lambda \in L^\bullet(\gamma)} W(\lambda)$$

with weights W defined in Definition 4.6.

Proof. The regular case of this theorem is proven in [12] and [14]. Graphs arise in the expression for $\omega_{g,n}$, as originally described in [17], as a means of encoding occurrences of the kernels $K(z_1, z_2)$ and $B(z_1, z_2)$ in formulae for the correlators obtained by iterating the recursion 1.2. Importantly the contributions to graphs by $B(p, p')$ are independent of the contributions to graphs by y . Hence the proof of Theorem 3 in [12] adapts immediately to allow irregular vertices, with changes only to vertex weights, and corresponding dilaton leaf weights, for vertices corresponding to irregular zeros of dx . The irregular vertex weights are determined from Proposition 4.2. \square

Remark 4.7. Although the main theorem in [12] requires a finiteness assumption on the spectral curve, Theorem 3 is proven there in full generality for any regular spectral curve. The theorem is also proven for any regular spectral curve in [14] Proposition 4.1, where the graphical sum is replaced by a moduli space of N -coloured stable Riemann surfaces.

The graphical expansion in Theorem 3 of the correlators $\omega_{g,n}$ of S can be restated in terms of differential operators acting on the partition function Z^S leading to a proof of the main result Theorem 1. The function $R(z)$ and operator \hat{R} below are obtained from [12].

Definition 4.8. Given a spectral curve $S = (\Sigma, B, x, y)$ define $R(z) = \sum R_k z^k \in \text{End}(V)[[z]]$ for a dimension N (= number of zeros of dx) vector space V by

$$\left[R^{-1}(z) \right]_{\beta}^{\alpha} = -\frac{\sqrt{z}}{\sqrt{2\pi}} \int_{\Gamma_{\beta}} B(\mathcal{P}_{\alpha}, p) \cdot e^{\frac{(x(\mathcal{P}_{\beta}) - x(p))}{z}}$$

and define the sequence $r_k \in \text{End}(V)$ by $R(z) = \exp\left(\sum_{\ell > 0} r_{\ell} z^{\ell}\right)$.

Definition 4.9. Given a spectral curve $S = (\Sigma, B, x, y)$ define

$$\begin{aligned} \hat{R} &= \exp \left\{ \sum_{\ell=1}^{\infty} \sum_{\alpha, \beta} \left(\sum_{k=0}^{\infty} v^{k, \beta} (r_k)_{\beta}^{\alpha} \frac{\partial}{\partial v^{k+\ell, \alpha}} + \frac{\hbar}{2} \sum_{m=0}^{\ell-1} (-1)^{m+1} (r_{\ell})_{\beta}^{\alpha} \frac{\partial^2}{\partial v^{m, \alpha} \partial v^{\ell-m-1, \beta}} \right) \right\} \\ \hat{T} &= \exp \left(\sum_{\alpha=1}^k \sum_{d \text{ odd}} d!! \frac{y_{d, \alpha}}{y_{-1, \alpha}} \frac{\partial}{\partial v^{d, \alpha}} + \sum_{\alpha=k+1}^N \sum_{d \text{ odd}} d!! \frac{y_{d, \alpha}}{y_{1, \alpha}} \frac{\partial}{\partial v^{d, \alpha}} \right) \\ \hat{\Delta} &\text{ acts on the } \alpha\text{th factor } Z^{\text{BGW}} \text{ by } \hbar \mapsto y_{-1, \alpha}^{-1} \hbar \text{ (or } y_{1, \alpha}^{-1} \hbar, \text{ for } Z^{\text{KW}}) \end{aligned}$$

We have assumed that y is irregular at \mathcal{P}_{α} for $\alpha = 1, \dots, k$ regular otherwise. Here r_k is defined in Definition 4.8 and $y_{2k+1, \alpha}$ are the odd coefficients of the local expansion $y = \sum y_{k, \alpha} z^k$ with respect to the local coordinate z defined by $x = \frac{1}{2}z^2 + x(\mathcal{P}_{\alpha})$.

Proof of Theorem 1. The graphical expansion in Theorem 3 corresponds to a Feynman expansion of the action of the exponential of a differential operator with quadratic and linear terms on a product of N functions of single variables. This viewpoint, i.e. the determination of the differential operators corresponding to the weighted graphical expressions is dealt with thoroughly in [13]. The weights give coefficients of differential operators. The factor $\prod_{v \in V(\gamma)} W(v)$ corresponds to the product $Z^{\text{BGW}}(\hbar, \{v^{d,1}\}) \dots Z^{\text{KW}}(\hbar, \{v^{d,N}\})$ of tau functions, the factor $\prod_{e \in E(\gamma)} W(e)$ corresponds to quadratic terms in the differential operator, the factor $\prod_{\ell \in L^*(\gamma)} W(\ell)$ corresponds to terms $t_j \frac{\partial}{\partial t_k}$ in the differential operator, and the factor $\prod_{\lambda \in L \bullet(\gamma)} W(\lambda)$ corresponds to translations.

It is crucial here that the definition of \hat{R} depends only on (Σ, B, x) which is independent of y and hence also of the property of the spectral curve being regular or irregular. This allows us to partly follow the proof of

the regular case in [12] since the same operator \hat{R} is being used. In [13] the equality

$$\begin{aligned} & \exp \left\{ \sum_{\ell=1}^{\infty} \sum_{\alpha, \beta} \left(\sum_{k=0}^{\infty} v^{k, \beta} (r_k)_{\beta}^{\alpha} \frac{\partial}{\partial v^{k+\ell, \alpha}} + \frac{\hbar}{2} \sum_{m=0}^{\ell-1} (-1)^{m+1} (r_{\ell})_{\beta}^{\alpha} \frac{\partial^2}{\partial v^{m, \alpha} \partial v^{\ell-m-1, \beta}} \right) \right\} \\ &= \exp \left\{ \sum_{\ell=1}^{\infty} \sum_{\alpha, \beta} \left(\sum_{k=0}^{\infty} v^{k, \beta} (r_k)_{\beta}^{\alpha} \frac{\partial}{\partial v^{k+\ell, \alpha}} \right) \right\} \exp \left\{ \sum_{k, \ell \geq 0} (2k-1)!! (2\ell-1)!! B_{2k, 2\ell}^{\alpha, \beta} \frac{\partial^2}{\partial v^{k, \alpha} \partial v^{\ell, \beta}} \right\} \end{aligned}$$

is proven via a variant of the Campbell–Baker–Hausdorff formula. Via the action of (constant coefficient) quadratic differential operators expressed graphically (also described well in [13]) this gives precisely the correct edge weights.

Remark 4.10. The general form of the Campbell–Baker–Hausdorff identity that we need pertains to three differential operators with constant coefficient matrices A , A' , and B (the matrices A and A' are symmetric): $A = (\vec{\partial} A \vec{\partial}^T)$, $A' = (\vec{\partial} A' \vec{\partial}^T)$, and $B = (\vec{t} B \vec{\partial}^T)$ with derivatives acting to the right. Then

$$e^{A+B} = e^{A'} e^B \text{ provided } B^T A' + A' B = e^{B^T} A e^B - A,$$

or

$$A' = \int_0^1 dx e^{xB^T} A e^{xB} \text{ or } A = \sum_{n=1}^{\infty} e^{-nB^T} (B^T A' + A' B) e^{-nB}.$$

The translation term follows immediately from Propositions 4.1, 4.2 and Corollary 4.3. This is a modification to the proof in [12] in two ways. Firstly, the proof in [12] uses an operator $\hat{R}_0 = \hat{R} \hat{T}_0$ where \hat{T}_0 is the translation arising out of a very special choice of y —see below for a discussion of this point. Secondly, here we needed Proposition 4.2 which resulted in terms given by translations of the tau function Z^{BGW} in addition to translations of the tau function Z^{KW} . \square

Note that related decompositions for matrix models appears in [1].

Example 4.11. The four decompositions of partition functions 1.1 follow from Theorem 1 applied to the following spectral curves. Each of the spectral curves is rational with common x and B given by the Cauchy kernel:

$$x = z + \frac{1}{z}, \quad B = \frac{dzdz'}{(z-z')^2}.$$

In particular, the operator \hat{R} is the same in all four examples. The examples differ by their choice of y :

- Z^{GW} uses $y = \ln z$ — [12, 29].
- Z^{GUE} uses $y = z$ — [16].
- Z^{Leg} uses $y = \frac{z}{z^2-1}$ — Section 5.
- Z^{Des} uses $y = \frac{z}{z+1}$ — [9].

Theorem 1 generalises a result for regular spectral curves in [12] and its proof is modeled on that result. We end this section with a comparison of the two decomposition formulae applied to regular spectral curves.

Definition 4.12. Given a regular spectral curve $S = (\Sigma, B, x, y)$ define Y_{α} locally in a neighbourhood of a zero \mathcal{P}_{α} of dx by the property that it satisfies

$$\frac{\sqrt{\zeta}}{\sqrt{2\pi}} \int_{\Gamma_{\alpha}} dY_{\alpha}(p) e^{(x(p)-x(\mathcal{P}_{\alpha}))\zeta} = \sum_{\beta=1}^N dy(\mathcal{P}_{\beta}) \cdot \frac{-1}{\sqrt{2\pi\zeta}} \int_{\Gamma_{\alpha}} B(p, \mathcal{P}_{\alpha}) e^{(x(p)-x(\mathcal{P}_{\alpha}))\zeta}.$$

The function Y_{α} is well-defined only up to addition of a function of x , because the Laplace transform annihilates functions of x , but this ambiguity disappears in the odd coefficients of the local expansion of Y_{α} which is all that is needed in the sequel.

The construction of [12] begins with Givental's decomposition of the partition function of (the correlators of) a cohomological field theory (with flat identity) and produces a spectral curve satisfying the restriction that $y = Y_\alpha$ for each $\alpha = 1, \dots, N$ (where $y = Y_\alpha$ up to local functions of x).

Theorem 4 ([12]). *Given a regular spectral curve $S = (\Sigma, B, x, y)$ satisfying $y = Y_\alpha$ for each $\alpha = 1, \dots, N$ the partition function Z^S satisfies the decomposition (1.3).*

The statement of this theorem is the reverse of the result in [12], but it is easily seen to be reversible on spectral curves satisfying the condition on y , [11]. The translation due to local expansions of y is implicit in the statement in [12] and appears inside an operator

$$\hat{R}_0 = \exp \left\{ - (r_k)_\mathbb{1}^\alpha \frac{\partial}{\partial v^{k+1, \alpha}} + \sum_{d=0}^{\infty} v^{d, \alpha} (r_k)_\alpha^\beta \frac{\partial}{\partial v^{k+d, \beta}} + \frac{\hbar}{2} \sum_{m=0}^{k-1} (-1)^{m+1} (r_k)^{\alpha, \beta} \frac{\partial^2}{\partial v^{m, \alpha} \partial v^{k-m-1, \beta}} \right\}$$

(which they denote by \hat{R}) where the vector $\mathbb{1} = \{dy(\mathcal{P}(\alpha))\}$. Raising of the indices of r_k is required for endomorphisms with respect to a general basis, but in this paper we express R in a basis (known as normalised canonical coordinates [21]) with respect to which the metric is the identity, so upper and lower indices are the same. It is related to the operator here by $\hat{R}_0 = \hat{R}\hat{T}$ for a translation \hat{T} .

This viewpoint allows us to interpret the decomposition (1.3) in the regular case as Givental's decomposition of the partition function of a cohomological field theory, now without the restriction of flat identity.

Corollary 4.13. *For any regular spectral curve $S = (\Sigma, B, x, y)$, Z^S is a partition function for a cohomological field theory.*

5. LEGENDRE ENSEMBLE

One application of Theorem 1 applied to the curve $(x^2 - 4)y^2 = 1$ is a Givental type decomposition for the partition function of the Legendre ensemble.

Proposition 5.1. *The resolvents for the Legendre ensemble*

$$(5.1) \quad \int_{H_N[-2,2]} \exp V(M) DM$$

with $V(M) = 0$ given by

$$W_g(x_1, \dots, x_n) = \left\langle \text{tr} \left(\frac{1}{x_1 - M} \right) \cdots \text{tr} \left(\frac{1}{x_n - M} \right) \right\rangle_g^c$$

satisfy topological recursion for the spectral curve

$$(5.2) \quad S = (\mathbb{P}^1, B = \frac{dzdz'}{(z-z')^2}, x = z + \frac{1}{z}, y = \frac{z}{z^2-1}).$$

In other words $W_g(x_1, \dots, x_n) dx_1 \dots dx_n$ gives an expansion of $\omega_{g,n}$ at $x_i = \infty$.

We omit the proof of Proposition 5.1 which is rather standard, via the usual loop equations.

A corollary of Theorem 1 is:

$$Z^{\text{Leg}}(\{v^{d,1}, v^{d,2}\}) = \hat{R}\hat{T}\hat{\Delta}Z^{\text{BGW}}(\{v^{d,1}\}) \times Z^{\text{BGW}}(\{v^{d,2}\}).$$

We now consider this example closely and demonstrate how the decomposition (1) produces calculations via pictures realising the graphical treatment of Theorem 3. We need only consider connected graphs to produce terms $F = \log Z$. In the calculations via pictures we instead use a related decomposition with only edge and vertex weights leading to simpler pictures. It is analogous to the decomposition proven by the first author in [5] and described next.

5.1. Chekhov-Givental decomposition. In [5] the first author obtained a decomposition for the Gaussian model that differs from (1.3). As in (1.3), it consists of the exponential of a quadratic differential operator acting on two copies of Z^{KW} , with the added translation term present in (1.3) but without a rotational term, whose absence was ensured by choosing global coordinates properly matching the local ones, see below. It is equivalent to (1.3) under a change of coordinates. For the Legendre hard-edge model we present here an analogous decomposition almost identical to the Gaussian model decomposition in [5]. As in the proof of Theorem 1, two ingredients of the decomposition arise from $B(p, q)$ – the Bergmann kernel – and from the 1-form ydx .

(i) ydx . We need ydx in global and local models.

(a) Global model: $x = e^\lambda + e^{-\lambda}$, $y = 1/(e^\lambda - e^{-\lambda})$, $ydx = d\lambda$

(b) Local models $y^2x = 1$, $x = \lambda^2$, $ydx = d\lambda$

So, global and local 1-forms coincide, $ydx^{\text{global}} = ydx^{\text{local}} = d\lambda$, which means that the linear differential part is absent in the Givental decomposition in this case.

(ii) $B(p, q)$. We again have global and local B's:

$$B(p, q)^{\text{global}} = \frac{de^\lambda de^\mu}{(e^\lambda - e^\mu)^2}, \quad B(p, q)^{\text{local}} = \frac{d\lambda d\mu}{(\lambda - \mu)^2},$$

and the corresponding integrals of B's, $e = \int^p \int^q B(\cdot, \cdot)$:

$$E(p, q)^{\text{global}} = \log(e^\lambda - e^\mu), \quad E(p, q)^{\text{local}} = \log(\lambda - \mu).$$

We then have that

$$\begin{aligned} B(p, q)^{\text{global}} &= \int_{C_{b_1}} \int_{C_{b_2}} B(p, \eta)^{\text{local}} E(\eta, \mu)^{\text{global}} B(q, \mu)^{\text{local}} \\ &= \delta_{b_1, b_2} B(p, q)^{\text{local}} + \int_{C_{b_1}} \int_{C_{b_2}} B(p, \eta)^{\text{local}} (E(\eta, \mu)^{\text{global}} - \delta_{b_1, b_2} E(\eta, \mu)^{\text{local}}) B(q, \mu)^{\text{local}}. \end{aligned}$$

In this expression b_1 and b_2 are the corresponding branch points, the first term in the right hand side describes the propagator of the local model, and the second term corresponds to the quadratic differential operator, which we consider in details below.

We interpret the propagators $B(p, \eta)^{\text{local}}$ and $B(q, \mu)^{\text{local}}$ as the correlation function endpoints;

the local times are $t_{2k+1}^- = \frac{1}{\lambda^{2k+1}}$ for $\lambda \sim 0$ and $t_{2k+1}^+ = \frac{1}{(\lambda - i\pi)^{2k+1}}$ for $\lambda \sim i\pi$; the corresponding global times are

$$(5.3) \quad t_{2k+1}^\pm = \frac{1}{(2k)!} \frac{\partial^{2k}}{\partial \lambda^{2k}} \frac{1}{\pm e^\lambda - 1}.$$

and the correlation function expansion is $W(\dots, x_j, \dots) \sim \sum_{r=1}^{\infty} \text{tr} H^r x_j^{-r-1}$ and recalling that

$$\langle \dots \text{tr}(H^r) \dots \rangle = r \frac{\partial}{\partial t_r} \langle \dots \rangle$$

we obtain that the expansion of the local correlation function has the form

$$W(\dots, x_j, \dots) = \sum_{r=1}^{\infty} r x_j^{-r-1} \frac{\partial}{\partial t_r}.$$

Here x_j are local coordinates (η and μ in our case).

When $b_1 = b_2$, the coefficient expansion of the quadratic differential operator stems from that of the function

$$E(\eta, \mu)^{\text{global}} - E(\eta, \mu)^{\text{local}} = \log \frac{e^\eta - e^\mu}{\eta - \mu} = \frac{\eta + \mu}{2} + \sum_{m=2}^{\infty} B_m \frac{(\eta - \mu)^m}{m \cdot m!},$$

where B_m are the Bernoulli numbers (recall that $B_{2s+1} = 0$ for $s \geq 1$), and with the accounting of the skew-symmetrisation under the transformations $\eta \leftrightarrow -\eta$ and $\mu \leftrightarrow -\mu$, we obtain the expansion

$$(5.4) \quad - \sum_{m=0}^{\infty} \frac{B_{2m+2}}{(2m+2)} \sum_{k=0}^m \frac{\eta^{2k+1}}{(2k+1)!} \frac{\mu^{2(m-k)+1}}{(2(m-k)+1)!}$$

with only odd powers of μ and η , so the corresponding differential operator, after evaluating the residues in η and μ , will contain derivatives only in odd times, as expected. This operator has the form:

$$(5.5) \quad A_{\pm, \pm} = -\frac{1}{2} \sum_{m=0}^{\infty} \frac{B_{2m+2}}{(2m+2)} \sum_{k=0}^m \frac{(2k+1)\partial/\partial t_{2k+1}^{\pm}}{(2k+1)!} \frac{(2(m-k)+1)\partial/\partial t_{2(m-k)+1}^{\pm}}{(2(m-k)+1)!}.$$

For different branching points, $b_1 = x^+$, $b_2 = x^-$, $E(\eta, \mu)^{\text{global}} = \log(e^\eta + e^\mu)$ and the corresponding expansion is

$$\log(e^\eta + e^\mu) = \frac{\eta + \mu}{2} + \sum_{m=2}^{\infty} \frac{B_m}{m} (1 - 2^m) \frac{(\eta - \mu)^m}{m!}$$

so after the skew-symmetrisation and substitution of differential operators in times, we obtain

$$(5.6) \quad A_{+, -} = - \sum_{m=0}^{\infty} \frac{B_{2m+2}}{(2m+2)} (1 - 2^{2m+2}) \sum_{k=0}^m \frac{(2k+1)\partial/\partial t_{2k+1}^+}{(2k+1)!} \frac{(2(m-k)+1)\partial/\partial t_{2(m-k)+1}^-}{(2(m-k)+1)!}.$$

The differential operators $A_{\pm, \pm}$ and $A_{+, -}$ exactly coincide with the corresponding operators from [5] where the analogous transformation was first derived in application to the Gaussian model.

Theorem 5. Consider the free energy of the Legendre ensemble on the interval $[-2, 2]$ with the potential $V(H) = \sum_{k=1}^{\infty} \frac{\tau_k}{k} H^k$ where the times τ_k are expressed using the Miwa-type transform $\tau_k = \sum_i (e^{\lambda_i} + e^{-\lambda_i})^{-k}$. The free energy of this model has the $1/N$ expansion of the form $F^{\text{Leg}} = \sum_{g=0}^{\infty} N^{2-2g} F_g^{\text{Leg}}$ where the term F_0^{Leg} is a quadratic polynomial in τ_k , and this term is to be interpreted as the normalisation. We then have the exact relation

$$(5.7) \quad e^{F^{\text{Leg}} - N^2 F_0^{\text{Leg}}} = e^{A_{+,+} + A_{-, -} + A_{+, -}} Z_{\text{BGW}}(t^+) Z_{\text{BGW}}(t^-),$$

where $Z_{\text{BGW}}(t^{\pm})$ are partition functions of the BGW model (exponentials of tau-functions of the KdV hierarchy) expressed in terms of times (5.3) and $A_{\cdot, \cdot}$ are quadratic differential operators (5.5) and (5.6).

Example 5.2. We now consider the decomposition for the F_2 term of the free energy.

$$F_2^{\text{Leg}} = \sum_{\pm} \text{circle}(2_{\pm}) + \sum_{\pm} \text{circle}(1_{\pm}) \text{---} \overset{A_{\pm, \pm}^{0,0}}{\text{circle}(1_{\pm})} + \text{circle}(1_{+}) \text{---} \overset{A_{+, -}^{0,0}}{\text{circle}(1_{-})} + \sum_{\pm} \text{circle}(1_{\pm}) \text{---} \overset{A_{\pm, \pm}^{0,0}}{\text{circle}(1_{\pm})}$$

Here

$$\text{circle}(2_{\pm}) = \frac{3\tau_1^{\pm}}{256(1 - \tau_0^{\pm})^3}, \quad \text{circle}(1_{\pm}) \text{---} = \frac{\partial F_1^{\text{Leg}}}{\partial \tau_0^{\pm}} = \frac{1}{8(1 - \tau_0^{\pm})}, \quad \text{circle}(1_{\pm}) \text{---} \overset{A_{\pm, \pm}^{0,0}}{\text{circle}(1_{\pm})} = \frac{\partial^2 F_1^{\text{Leg}}}{\partial (\tau_0^{\pm})^2} = \frac{1}{8(1 - \tau_0^{\pm})^2},$$

$$A_{\pm, \pm}^{k,l} = -\frac{1}{2} \cdot \frac{B_{2(k+l+1)}}{2(k+l+1)(2k)!(2l)!}, \quad A_{+, -}^{k,l} = -\frac{B_{2(k+l+1)}}{2(k+l+1)(2k)!(2l)!} (2^{2(k+l+1)} - 1).$$

Recalling that $B_2 = 1/6$, we have that $A_{\pm,\pm}^{0,0} = -1/24$ and $A_{+,-}^{0,0} = -1/4$, so the combination of the second and fourth terms in the graphical expansion for F_2^{Leg} yield

$$\left(-\frac{1}{8^2 \cdot 24} - \frac{1}{8 \cdot 24}\right) \frac{1}{(1 - \tau_0^\pm)^2} = -\frac{3}{512} \frac{1}{(1 - \tau_0^\pm)^2}$$

and the third term of the same graphical expansion yields

$$-\frac{1}{8^2 \cdot 4} \frac{1}{(1 - \tau_0^+)(1 - \tau_0^-)} = -\frac{1}{256} \frac{1}{(1 - \tau_0^+)(1 - \tau_0^-)}$$

in a perfect agreement with the algebraic expression for F_2^{Leg} given in the exact formulae below. We reproduce all terms in the second line of this expression. The whole first line, as was demonstrated above, just combines into $\frac{3\tau_1^\pm}{256(1 - \tau_0^\pm)^3}$ (recall that normalisations for t_k and τ_k differ by the factor of 4^k).

Example 5.3. We now consider the decomposition for F_3 term of the free energy.

$$\begin{aligned}
F_3^{\text{Leg}} = & \sum_{\pm} \text{[Diagram: circle with } 3_{\pm} \text{]} + \sum_{\pm} 2 \text{ [Diagram: } 2_{\pm} \text{ --- } 1_{\pm} \text{]}^{A_{\pm,\pm}^{0,0}} + 1 \text{ [Diagram: } 2_{\pm} \text{ --- } 1_{\mp} \text{]}^{A_{+,-}^{0,0}} + \sum_{\pm} 1 \text{ [Diagram: } 2_{\pm} \text{ with dashed circle]}^{A_{\pm,\pm}^{0,0}} \\
& + \sum_{\pm} 2 \times 2 \text{ [Diagram: } 2_{\pm} 1 \text{ --- } 1_{\pm} \text{]}^{A_{\pm,\pm}^{1,0}} + 1 \times 2 \text{ [Diagram: } 2_{\pm} 1 \text{ --- } 1_{\mp} \text{]}^{A_{\pm,\mp}^{1,0}} + \sum_{\pm} 2 \times 2 \text{ [Diagram: } 2_{\pm} \text{ with dashed circle]}^{A_{\pm,\pm}^{1,0}} \\
& + \sum_{\pm} 1/2 \text{ [Diagram: } 1_{\pm} \text{ with dashed circle]}^{A_{\pm,\pm}^{0,0}} + \sum_{\pm} 2 \text{ [Diagram: } 1_{\pm} \text{ --- } 1_{\pm} \text{]}^{A_{\pm,\pm}^{0,0}} + \sum_{\pm} 1 \text{ [Diagram: } 1_{\pm} \text{ --- } 1_{\mp} \text{]}^{A_{\pm,\mp}^{0,0}} \\
& + \sum_{\pm} 1 \text{ [Diagram: } 1_{\pm} \text{ --- } 1_{\pm} \text{]}^{A_{\pm,\pm}^{0,0}} + 1/2 \text{ [Diagram: } 1_{+} \text{ --- } 1_{-} \text{]}^{A_{+,-}^{0,0}} + \sum_{\pm} 2 \text{ [Diagram: } 1_{\pm} \text{ --- } 1_{\pm} \text{]}^{A_{\pm,\pm}^{0,0}} + \sum_{\pm} 2 \text{ [Diagram: } 1_{\pm} \text{ --- } 1_{\pm} \text{]}^{A_{\pm,\pm}^{0,0}} \\
& + \sum_{\pm} 2 \text{ [Diagram: } 1_{\pm} \text{ --- } 1_{\pm} \text{]}^{A_{\pm,\pm}^{0,0}} + \sum_{\pm} 1/2 \text{ [Diagram: } 1_{\pm} \text{ --- } 1_{\mp} \text{]}^{A_{\pm,\mp}^{0,0}} + \sum_{\pm} 1/2 \text{ [Diagram: } 1_{\pm} \text{ --- } 1_{\mp} \text{]}^{A_{\mp,\pm}^{0,0}}
\end{aligned}$$

Here

$$\text{[Diagram: } 2_{\pm} 1 \text{]} = \frac{3}{256(1 - \tau_0^\pm)^3}, \quad \text{[Diagram: } 2_{\pm} 1 \text{]} = \frac{3 \cdot 3}{256(1 - \tau_0^\pm)^4}, \quad \text{[Diagram: } 2_{\pm} \text{ with dashed circle]} = \frac{3 \cdot 3 \cdot 4\tau_1^\pm}{256(1 - \tau_0^\pm)^5},$$

and

$$\text{[Diagram: } 1_{\pm} \text{ with dashed circle]}^k = \frac{\partial^k F_1^{\text{Leg}}}{\partial (\tau_0^\pm)^k} = \frac{(k-1)!}{8(1 - \tau_0^\pm)^k},$$

$$A_{\pm, \pm}^{k,l} = -\frac{1}{2} \cdot \frac{B_{2(k+l+1)}}{2(k+l+1)(2k)!(2l)!}, \quad A_{+,-}^{k,l} = -\frac{B_{2(k+l+1)}}{2(k+l+1)(2k)!(2l)!} (2^{2(k+l+1)} - 1).$$

$F_3^{\text{BGW}} = C_1 \tau_2 / (1 - \tau_0)^6 + C_2 (\tau_1)^2 / (1 - \tau_0)^7$. Here τ_2^\pm is obtained by a Laplace transform from a “pure” state corresponding to the monomial $b^4(\mp 1)^b$. Recalling that $B_2 = 1/6$ and $B_4 = -1/30$, we have that $A_{\pm, \pm}^{0,0} = -1/24$, $A_{+,-}^{0,0} = -1/4$, $A_{\pm, \pm}^{1,0} = 1/(30 \cdot 16)$, $A_{\pm, \mp}^{1,0} = 1/16$. so we fix everything indicating in gray the symmetry factors of diagrams. All derivatives in the times τ_1^\pm are to be multiplied by factors of two in order to obtain the known answer for F_3 given in the exact formulae below.

The terms in the first line are

$$F_3^{\text{BGW}}(\tau_+) + F_3^{\text{BGW}}(\tau_-) - \frac{2 \cdot 3 \cdot 3 \cdot \tau_1^\pm}{24 \cdot 256 \cdot 8(1 - \tau_0^\pm)^5} - \frac{3 \cdot 3 \cdot \tau_1^\pm}{4 \cdot 256 \cdot 8(1 - \tau_0^\pm)^4(1 - \tau_0^\mp)} - \frac{3 \cdot 3 \cdot 4 \cdot \tau_1^\pm}{24 \cdot 256(1 - \tau_0^\pm)^5}$$

and all other terms depend only on zeroth times and are a linear combination of $(1 - \tau_0^+)^{-4} + (1 - \tau_0^-)^{-4}$, $(1 - \tau_0^+)^{-3}(1 - \tau_0^-)^{-1} + (1 - \tau_0^-)^{-3}(1 - \tau_0^+)^{-1}$, and $(1 - \tau_0^+)^{-2}(1 - \tau_0^-)^{-2}$ with necessarily positive coefficients.

$$\tau_1^\pm = 4t_1^1 \mp 4t_1^0 \mp t_0^1,$$

$$\tau_2^\pm = 32t_2^0 + 16t_1^1 \mp 32t_2^1 \mp 8t_1^0 \mp t_0^1,$$

Unlike the Gaussian case, we see that the free energy F_g^{Leg} is in some sense finite. It is a finite sum of rational functions apparent in the exact formulae below.

Exact formulae.

$$F_1^{\text{Leg}} = \frac{1}{8}(1 - \log(1 - t_0^0 - t_0^1)) + \frac{1}{8}(1 - \log(1 - t_0^0 + t_0^1))$$

$$F_2^{\text{Leg}} = \frac{3}{64} \frac{t_0^1 + t_1^1}{(1 - t_0^0 - t_0^1)^3} + \frac{3}{64} \frac{t_0^1 - t_1^1}{(1 - t_0^0 + t_0^1)^3} + \frac{3}{256} \frac{t_0^1}{(1 - t_0^0 - t_0^1)^3} - \frac{3}{256} \frac{t_0^1}{(1 - t_0^0 + t_0^1)^3} \\ - \frac{3}{512} \frac{1}{(1 - t_0^0 - t_0^1)^2} - \frac{3}{512} \frac{1}{(1 - t_0^0 + t_0^1)^2} - \frac{1}{256} \frac{1}{(1 - t_0^0 - t_0^1)(1 - t_0^0 + t_0^1)}$$

$$F_3^{\text{Leg}} = \frac{15 t_2^0 + 15 t_2^1}{256 (1 - t_0^0 - t_0^1)^5} + \frac{15 t_2^0 - 15 t_2^1}{256 (1 - t_0^0 + t_0^1)^5} + \frac{63 (t_1^0 + t_1^1)^2}{256 (1 - t_0^0 - t_0^1)^6} + \frac{63 (-t_1^0 + t_1^1)^2}{256 (1 - t_0^0 + t_0^1)^6} \\ + \frac{1}{(1 - t_0^0 - t_0^1)^5} \left(-\frac{21 t_1^0}{2048} + \frac{9 t_1^1}{2048} \right) + \frac{1}{(1 - t_0^0 + t_0^1)^5} \left(\frac{21 t_1^0}{2048} + \frac{9 t_1^1}{2048} \right) \\ + \frac{1}{(1 - t_0^0 - t_0^1)^4 (1 - t_0^0 + t_0^1)} \left(-\frac{9 t_1^0}{2048} - \frac{9 t_1^1}{2048} \right) + \frac{1}{(1 - t_0^0 + t_0^1)^4 (1 - t_0^0 - t_0^1)} \left(\frac{9 t_1^0}{2048} - \frac{9 t_1^1}{2048} \right) \\ + \frac{1}{(1 - t_0^0 - t_0^1)^6} \left(\frac{63 t_1^0 t_0^1}{512} + \frac{63 t_1^1 t_0^1}{512} \right) + \frac{1}{(1 - t_0^0 + t_0^1)^6} \left(\frac{63 t_1^0 t_0^1}{512} - \frac{63 t_1^1 t_0^1}{512} \right) \\ + \frac{63 t_0^1^2}{4096 (1 - t_0^0 - t_0^1)^6} + \frac{63 t_0^1^2}{4096 (1 - t_0^0 + t_0^1)^6} - \frac{9 t_0^1}{2048 (1 - t_0^0 - t_0^1)^5} + \frac{9 t_0^1}{2048 (1 - t_0^0 + t_0^1)^5} \\ + \frac{9}{16384 (1 - t_0^0 - t_0^1)^4} + \frac{9}{16384 (1 - t_0^0 + t_0^1)^4} + \frac{9}{8192 (1 - t_0^0 - t_0^1)^3 (1 - t_0^0 + t_0^1)} \\ + \frac{9}{8192 (1 - t_0^0 - t_0^1) (1 - t_0^0 + t_0^1)^3} + \frac{5}{8192 (1 - t_0^0 + t_0^1)^2 (1 - t_0^0 - t_0^1)^2}$$

Furthermore, $F_g^{\text{Leg}}|_{t_0^0=0=t_1^0}$ is a polynomial, in particular it has finitely many terms. Moreover, $F_g^{\text{Leg}}|_{t_0^0=0=t_1^0}$ uniquely determines F_g^{Leg} via the so-called dilaton and divisor equations.

The dilaton equation reflects the fact F_g is homogeneous of degree $2 - 2g$ (if we put $q = t_0^0 - 1$, known as a dilaton shift). This is clear in the examples above. It is easily obtained via loop equations for the matrix integral, and is a consequence of topological recursion.

$$(5.8) \quad \sum_{z=\alpha} \text{Res} \phi(z) \omega_{g,n+1}(z, z_1, \dots, z_n) = (2g - 2 + n) \omega_{g,n}(z_1, \dots, z_n)$$

where $d\phi = ydx = dz/z$ hence $\phi = \log z$. Equivalently

$$(5.9) \quad \frac{\partial}{\partial t_0^0} F_g = \sum_k \left\{ (2k+1)t_k^0 \frac{\partial}{\partial t_k^0} + (2k+1)t_k^1 \frac{\partial}{\partial t_k^1} + (2g-2) \right\} F_g.$$

The divisor equation is subtler. It is obtained using:

$$(5.10) \quad \omega_{g,n}(z_1, z_2, \dots, z_n) = \sum_{z=\alpha} \text{Res} K(z_1, z) \left[\omega_{g-1, n+1}(z, \hat{z}, z_2, \dots, z_n) + \sum_{\substack{g_1+g_2=g \\ I \sqcup J = \{2, \dots, n\}}} \omega_{g_1, |I|+1}(z, z_1) \omega_{g_2, |J|+1}(\hat{z}, z_j) \right]$$

to get:

$$\text{Res}_{z_1=0} x(z_1) \omega_{g,n}(z_1, z_2, \dots, z_n) = \text{Res}_{z_1=0} x(z_1) \sum_{z=\alpha} \text{Res} K(z_1, z) \sum_{j=2}^n \omega_{g,n-1}(z, z_2, \dots, \hat{z}_j, \dots, z_n) \omega_{0,2}(\hat{z}, z_j)$$

which becomes

$$(5.11) \quad \frac{\partial}{\partial t_0^1} F_g = \sum_k \left\{ \left[(2k+1)t_k^1 + \binom{k}{2} t_{k-1}^0 \right] \frac{\partial}{\partial t_k^0} + \left[(2k+1)t_k^0 + \binom{k+1}{2} t_{k-1}^1 \right] \frac{\partial}{\partial t_k^1} \right\} F_g.$$

The dilaton and divisor equations uniquely determine F_g^{Leg} from the initial conditions $F_g^{\text{Leg}}|_{t_0^0=0=t_1^0}$ given by finitely many terms.

APPENDIX A QUANTISATION

In this section we give a brief background for the construction of \hat{R} from R following Givental [20]. We first consider quantisation in finite dimensions which easily generalises to infinite dimensions.

Consider the standard holomorphic form on $\mathbb{C}^{2N} = T^*\mathbb{C}^N$ given by $\omega = \sum dp_\alpha \wedge dq_\alpha$ with Darboux coordinates $\{q_\alpha, p_\alpha\}$. A transformation $A : \mathbb{C}^{2N} \rightarrow \mathbb{C}^{2N}$ that preserves the symplectic form is *symplectic*. We will consider only linear symplectic transformations which correspond to matrices $A \in \text{Sp}(2N, \mathbb{C})$. So-called *infinitesimal* symplectic transformations, corresponding to elements of the Lie algebra $\text{sp}(2N, \mathbb{C})$, give rise to vector fields that preserve ω , known as Hamiltonian vector fields. The $2N^2 + N$ -dimensional space $\text{sp}(2N, \mathbb{C})$ is isomorphic to the $2N^2 + N$ -dimensional vector space of Hamiltonians

$$p_\alpha p_\beta, \quad p_\alpha q_\beta, \quad q_\alpha q_\beta$$

so $H(p, q) = p_1 q_2$ for example. Quantisation of these coordinates, i.e. promotion to operators $p_\alpha \mapsto \hat{p}_\alpha$ and $q_\alpha \mapsto \hat{q}_\alpha$, satisfies

$$[\hat{p}_\alpha, \hat{q}_\beta] = \delta_{\alpha\beta} \hbar, \quad [\hat{p}_\alpha, \hat{p}_\beta] = 0 = [\hat{q}_\alpha, \hat{q}_\beta]$$

so we naturally choose $\hat{p}_\alpha = \hbar \frac{\partial}{\partial q_\alpha}$ and \hat{q}_α acts by multiplication by q_α . Quantisation of a function, or observable, in $\{p_\alpha, q_\beta\}$ is not unique. We can consistently define quantisation of the quadratic Hamiltonians by

$$\widehat{p_\alpha p_\beta} = \hbar^2 \frac{\partial^2}{\partial q_\alpha \partial q_\beta}, \quad \widehat{p_\alpha q_\beta} = \hbar q_\beta \frac{\partial}{\partial q_\alpha}, \quad \widehat{q_\alpha q_\beta} = q_\alpha q_\beta.$$

Linear combinations of these give the quantisation of infinitesimal symplectic transformations. Hence we can define the quantisation of a linear symplectic transformation $A = \exp a$ for $a \in \mathfrak{sp}(2N, \mathbb{C})$ to be $\hat{A} = \exp \hat{a}$.

This construction generalises to the infinite dimensional symplectic manifold $\mathcal{H} = H[[z, z^{-1}]]$ for $H \cong \mathbb{C}^N$ with symplectic form

$$\Omega(f(z), g(z)) = \operatorname{Res}_{z=0} f(-z)g(z)dz.$$

$\mathcal{H}_+ = H[[z]]$ is Lagrangian with respect to Ω and $(\mathcal{H}, \Omega) \cong (T^*\mathcal{H}_+, \omega^{\text{canonical}})$. Darboux coordinates for Ω are $q_{k,\alpha}, p_{k,\alpha}$ defined by $\mathcal{H} \ni f(z) = \sum_{k \geq 0} q_{k,\alpha} z^k + \sum_{k < 0} p_{k,\alpha} z^{-k}$.

Given a dimension N vector space V and a sequence of operators $r_k : V \rightarrow V, k = 1, 2, \dots$ such that $r_k(-v) = (-1)^{k+1} r_k(v)$ define

$$\widehat{r_k z^k} := \sum_{d=0}^{\infty} v^{d,\alpha} (r_k)_\alpha^\beta \frac{\partial}{\partial v^{k+d,\beta}} + \frac{\hbar}{2} \sum_{m=0}^{k-1} (-1)^{m+1} (r_k)^{\alpha,\beta} \frac{\partial^2}{\partial v^{m,\alpha} \partial v^{k-m-1,\beta}}.$$

Then the quantisation of $R(z) = \exp r(z)$ is given by $\hat{R}(z) = \exp \hat{r}(z)$.

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STEKLOV MATHEMATICAL INSTITUTE AND LAB. PONCELET, MOSCOW, RUSSIA; NIELS BOHR INSTITUTE, COPENHAGEN, DENMARK

SCHOOL OF MATHEMATICS AND STATISTICS, UNIVERSITY OF MELBOURNE, VIC 3010 AUSTRALIA

E-mail address: chekhov@mi.ras.ru, pnorbury@ms.unimelb.edu.au